

## RECENT IMPROVEMENTS IN MIXED/ENHANCED SHELL ELEMENTS WITH DRILLING ROTATION

K. Wiśniewski<sup>1</sup>, E. Turska<sup>2</sup>

<sup>1</sup>Institute of Fundamental Technological Research, Warsaw, Poland

<sup>2</sup>Polish Japanese Institute of Information Technology, Warsaw, Poland

### 1. Introduction

The purpose of the presentation is to summarize features of our recent formulation of mixed/enhanced four-node shell elements based on the Hu-Washizu (HW) functional *with* rotational degrees of freedom and illustrate their numerical performance [1–5].

The drilling rotation, defined as the rotation about the shell director, is not naturally present in the shell equations derived from the non-polar Cauchy continuum for the Reissner hypothesis and Green strain. To alleviate this restriction, we consider the *extended configuration* space, consisting of the deformation  $\chi \in R^3$  and rotations  $\mathbf{Q} \in SO(3)$  subjected to the Rotation Constraint (RC) equation  $skew(\mathbf{Q}^T \mathbf{F}) = \mathbf{0}$ , where  $\mathbf{F} = \nabla \chi$ . For shells,  $\mathbf{Q}$  is reduced to the rotation at the reference surface  $\mathbf{Q}_0$  and only the drilling component (12) of the RC equation,  $[skew(\mathbf{Q}_0^T \mathbf{F})]_{12} = 0$ , is used. This approach enables us to use standard constitutive equations for shells; in contrast to the Cosserat-type shells, which also use the drilling rotation but need specialized constitutive equations.

### 2. Shell pure HW functional with rotations

Mixed finite elements were pioneered by Pian in 1964 and a lot of work has been done since then to improve their theoretical foundations and performance. In the class of four-node elements, the best seem to be the ones based on the three-field Hu-Washizu (HW) functional; they exhibit a higher accuracy of displacements and stresses and significantly better convergence properties in non-linear problems than other elements, such as the enhanced strain (EAS or EADG) elements.

Our formulation is based on the 2nd Piola-Kirchhoff stress  $\mathbf{S}$  and the Green strain  $\mathbf{E}(\mathbf{u}) = \frac{1}{2}(\mathbf{F}^T \mathbf{F} - \mathbf{I})$ . For 3D formulation, the Lagrange multiplier method is used to append the RC equation to the HW functional,

$$(1) F_{\text{HW}}(\mathbf{u}, \mathbf{Q}, \mathbf{S}^*, \mathbf{E}^*, \mathbf{T}^*) = \int_{\mathbf{v}} \{W(\mathbf{E}^*) + \mathbf{S}^* \cdot [\mathbf{E}(\mathbf{u}) - \mathbf{E}^*] + \mathbf{T}^* \cdot skew(\mathbf{Q}^T \mathbf{F})\} dV - F_{\text{ext}},$$

where  $W(\mathbf{E}^*)$  is the strain energy expressed by the independent strain  $\mathbf{E}^*$ , and  $\mathbf{S}^*$  is the independent stress which serves as the Lagrange multiplier for the difference of strains  $\mathbf{E}^*$  and  $\mathbf{E}(\mathbf{u})$ . Besides,  $\mathbf{T}^*$  is the skew-symmetric Lagrange multiplier for the RC equation  $skew(\mathbf{Q}^T \mathbf{F}) = \mathbf{0}$ .

To derive the HW functional for shells, we use the Reissner shell kinematics,

$$(2) \quad \mathbf{x}(\zeta) = \mathbf{x}_0 + \zeta \mathbf{Q}_0 \mathbf{t}_3, \quad \zeta \in [-1, +1],$$

where  $\mathbf{x}_0$  is a position of the reference surface and  $\mathbf{t}_3$  is the shell director. The rotation of the reference surface  $\mathbf{Q}_0$  is parameterized by the canonical rotation vector. For this kinematics, the Green strain linearized in  $\zeta$  becomes  $\mathbf{E}(\zeta) \approx \boldsymbol{\varepsilon} + \zeta \boldsymbol{\kappa}$ . Let us also assume

the independent strain in the analogous form,  $\mathbf{E}^*(\zeta) \approx \boldsymbol{\varepsilon}^* + \zeta \boldsymbol{\kappa}^*$ . Then, by integration of the HW functional of Eq. (1) over the shell thickness, we obtain its shell counterpart,

$$(3) \quad F_{\text{HW}}^{\text{sh}} = \int_{\mathbf{A}} \{W^{\text{sh}}(\boldsymbol{\varepsilon}^*, \boldsymbol{\kappa}^*) + \mathbf{N}^* \cdot [\boldsymbol{\varepsilon}(\mathbf{u}, \mathbf{Q}_0) - \boldsymbol{\varepsilon}^*] + \mathbf{M}^* \cdot [\boldsymbol{\kappa}(\mathbf{u}, \mathbf{Q}_0) - \boldsymbol{\kappa}^*] + F_{\text{RC}}^{\text{drill}}\} dA - F_{\text{ext}}^{\text{sh}},$$

where  $\mathbf{N}^*$  and  $\mathbf{M}^*$  are defined as the integrals of  $\mathbf{S}^*$  and  $\zeta \mathbf{S}^*$  over the shell thickness, and  $\mathbf{A}$  is the area of the shell reference surface. The drilling RC part is used in the Perturbed Lagrange (PL) form,

$$(4) \quad F_{\text{RC}}^{\text{drill}} = T^* [\text{skew}(\mathbf{Q}_0^{\text{T}} \mathbf{F})]_{12} - (1/2\gamma)(T^*)^2$$

where  $T^*$  is the (scalar) Lagrange multiplier and  $\gamma \in (0, \infty)$  is the regularization parameter. Note that seven fields ( $\mathbf{u}, \mathbf{Q}_0, \mathbf{N}^*, \mathbf{M}^*, \boldsymbol{\varepsilon}^*, \boldsymbol{\kappa}^*, T^*$ ) are used in  $F_{\text{HW}}^{\text{sh}}$ . This methodology yields pure shell HW functional, the shell element for which performs very well but requires a large number of parameters and so is not efficient.

### 3. Shell partial HW functionals with rotations

In an alternative approach, we derive the shell HW functional differently. First, the 3D potential energy functional is integrated, to obtain its shell counterpart

$$(5) \quad F_{\text{PE}}^{\text{sh}} = \int_{\mathbf{A}} W^{\text{sh}}(\boldsymbol{\varepsilon}, \boldsymbol{\kappa}) dA - F_{\text{ext}}^{\text{sh}}.$$

Then, we construct the *partial* shell HW functional for selected strain components only while still using the potential energy functional for the other components. The so-derived functionals can be used to select the “optimal” formulation, which ensures the required numerical performance using a minimum number of parameters. For instance, in the HW29 shell element, we use the HW functional only for the 0-th order strain components  $\varepsilon_{\alpha\beta}$  and  $\varepsilon_{\alpha 3}$  ( $\alpha, \beta = 1, 2$ ),

$$(6) \quad F_{\text{HW}}^{\text{sh}} = \int_{\mathbf{A}} \{W^{\text{sh}}(\varepsilon_{\alpha\beta}^*, \varepsilon_{\alpha 3}^*, \kappa_{\alpha\beta}) + N_{\alpha\beta}^* [\varepsilon_{\alpha\beta} - \varepsilon_{\alpha\beta}^*] + N_{\alpha 3}^* [\varepsilon_{\alpha 3} - \varepsilon_{\alpha 3}^*] + F_{\text{RC}}^{\text{drill}}\} dA - F_{\text{ext}}^{\text{sh}}.$$

Using such a methodology, we have developed several mixed/enhanced shell elements based on the Hu-Washizu functional *with* rotational degrees of freedom, [1–4]. Numerical examples will illustrate various aspects of their performance, such as: accuracy, radius of convergence, required number of iterations of the Newton method or the arc-length method and time of computations. Additionally, examples enabling comparisons with our new HW elements using only translational degrees of freedom will be provided.

### References

1. K. Wiśniewski, E. Turska (2012). Four-node mixed Hu-Washizu shell element with drilling rotation, *Int. J. Num. Meth. Engng*, **90**, 506–536.
2. K. Wiśniewski (2010). *Finite Rotation Shells. Basic Equations and Finite Elements for Reissner Kinematics*, CIMNE-Springer.
3. K. Wiśniewski, W. Wagner, E. Turska, F. Gruttmann (2010). Four-node Hu-Washizu elements based on skew coordinates and contravariant assumed strain, *Computers & Structures*, **88**, 1278–1284.
4. K. Wiśniewski, E. Turska (2009). Improved four-node Hu-Washizu elements based on skew coordinates, *Computers & Structures*, **87**, 407–424.
5. K. Wiśniewski, E. Turska (2006). Enhanced Allman quadrilateral for finite drilling rotations, *Comp. Meth. Appl. Mech. Engng.*, **195**, 44–47, 6086–6109.